

# Aarhus Master Class

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$X$  a compact riemannian (usually) 4-manifold  $X$  possibly with boundary.

$G \subset \text{Aut}(V)$  a compact Lie group.

$\mathfrak{g} \subset \text{End}(V)$ , Lie algebra of  $G$ .

$\pi : P \rightarrow X$  a principal  $G$  bundle.  $E = P \times_G V$  denote the associated  $V$  bundle. The vertical tangent space

$VT_p P = \ker d\pi$ . There is a canonical isomorphism of  $\iota_p : VT_p P \rightarrow \mathfrak{g}$ .

$\text{ad}P = P \times_{\text{ad}} \mathfrak{g} \subset \text{End}(E)$  and  $\text{Ad}P = P \times_{\text{Ad}} G \subset \text{Aut}(E)$ .  
Forms with values in  $\text{ad}P$  are a super lie algebra.

$$[\cdot \wedge \cdot] : \Lambda^i(T^*X) \otimes \text{ad}P \otimes \Lambda^j(T^*X) \otimes \text{ad}P \rightarrow \Lambda^{i+j}(T^*X) \otimes \text{ad}P$$

so that

$$[a \wedge b] = (-1)^{|a|+|b|+1} [b \wedge a]$$

For us  $G$  will be  $SO_3$  or  $SU_2$ .

A **connection**  $A$  in  $P$  can be viewed many ways.

- A system of parallel transport in  $P$

The **curvature** of a connection measures the failure of  $H_A$  to be involutive. As a two form on  $P$

$$F_A = dA + \frac{1}{2}[A \wedge A].$$

Can view  $F_A \in C^\infty(X, \Lambda^2 \otimes \text{ad}P)$ .

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- A covariant derivative

$$d_A : C^\infty(X; \text{ad}P) \rightarrow C^\infty(X, T^*X \otimes \text{ad}P).$$

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The Yang-Mills functional. The Yang-Mills energy of  $A$  is

$$E(A) = - \int_X \operatorname{tr}(F_A \wedge *F_A) = \int_X |F_A|^2 * 1.$$

(Recall the  $\operatorname{tr}(A^2)$  is a negative definite form on **skew symmetric** matrices.) The Euler-Lagrange equations for  $E$  is

$$d_A^* F_A = 0$$

the **Yang-Mills Equations**. Recall that the curvature always satisfies the Bianchi Identity.

$$d_A F_A = 0$$

Informally  $A$  is a non-linear harmonic form!

The Chern-Weil formula. Consider:

$$\int_X \text{tr}(F_A \wedge F_A).$$

If  $X$  is closed then this integral depends only on  $P$  and not on  $A$ .

Proof:  $A, A' = A + a$ . Let  $A_t = A + ta$ . Then

$$F_{A_t} = F_A + td_A a + \frac{t^2}{2}[a \wedge a].$$

$$\begin{aligned} \frac{d}{dt} \int_X \text{tr}(F_{A_t} \wedge F_{A_t}) &= 2 \int_X \text{tr}(d_{A_t} a \wedge F_{A_t}) \\ &= 2 \int_X d \text{tr}(a \wedge F_{A_t}) \\ &= 0. \end{aligned}$$



If  $G = SO_3$  then we define the instanton number:

$$k = -\frac{1}{4} \langle p_1(P), [X] \rangle = -\frac{1}{8\pi^2} \int_X \text{tr}(F_A \wedge F_A) \in \frac{1}{4}\mathbb{Z}$$

Note if  $P$  lifts to an  $SU_2$  bundle  $Q$  then  $\text{ad}Q \equiv \text{ad}P$  and so

$$\begin{aligned} p_1(P) &= p_1(\text{ad}P) = p_1(\text{ad}Q) \\ &= -c_2(\text{ad}Q \otimes \mathbb{C}) = -c_2(\text{End}(E)) = -4c_2(E) \end{aligned}$$

where  $E$  denotes the  $\mathbb{C}^2$ -bundle associated to  $Q$ . and similarly if  $P$  lifts to a  $U_2$  bundle  $Q$

$$p_1(P) = -4c_2(Q) + c_1^2(Q).$$

$P$  is determined upto isomorphism by  $p_1(P)$  and  $w = w_2(P) \in H^2(X, \mathbb{Z}_2)$ .

## Spaces of connections and gauge transformations.

$A$  a connection.  $\mathcal{A}$  is the space of  $C^\infty$ -connections.  $\mathcal{A}$  is an affine space for the space

$$\Omega^1(X; \text{ad}P) = C^\infty(X; T^*X \otimes \text{ad}P).$$

$g : P \rightarrow P$  is an automorphism (**gauge transformation**) of  $P$ .

$g$  is a section of the bundle  $\text{Ad}P$ .

$\mathcal{G}$  the space of  $C^\infty$ -sections of  $\text{Ad}P$  or **group of gauge transformations**.

$\nabla_A f$  or  $d_A f$  The induced covariant derivative in any associated bundle,  $F = P \times_\rho V$ .

The action of  $g \in \mathcal{G}$ .

$$g \cdot A = A + g d_A g^{-1}.$$

$d_A g^{-1}$  means that covariant derivative of  $g^{-1}$  thought of as a section of  $\text{End}(E) = E^* \otimes E$  induced by  $A$ . Locally

We are interested in the studying connections **up to isomorphism** i.e.

$$B = \mathcal{A}/\mathcal{G}.$$

We need to understand how bad the action is.

### Lemma

*$g \cdot A = A$  if and only if  $g$  is a parallel section of  $\text{End}(E)$ .*

Proof.

$$g \cdot A = A \implies g d_A g^{-1} = -(d_A g) g^{-1} = 0.$$

and hence  $d_A g = 0$ .  $\square$

**N.B** When  $G = SO_3$  the we can form the bundle  $\text{Ad}^1 P$  of "determinant 1" gauge transformations

$$\text{Ad}^1 P = P \times_{SO_3} SU_2.$$

where  $SO_3 = SU_2/Z(SU_2)$  acts on  $SU_2$  by conjugation. The sections of this bundle  $\mathcal{G}^1$  maps to  $\mathcal{G}$ .

### Exercise

*The map  $\mathcal{G}^1 \rightarrow \mathcal{G}$  has kernel  $\mathbb{Z}_2$  if  $X$  is connected. The cokernel isomorphic to  $H^1(X; \mathbb{Z}_2)$ .*

We will often want to work with  $\mathcal{A}/\mathcal{G}^1$  rather than  $\mathcal{A}/\mathcal{G}$ . To deal with later we'll choose submanifold Poincaré dual to  $w_2(P)$ .

Given a loop  $\gamma$  based at  $x_0 \in X$  the **holonomy** of  $A$  along  $\gamma$  is the automorphism of  $P|_{x_0}$  given by parallel transport along  $\gamma$ .

$$\text{hol}_A(\gamma) \in \text{Ad}P|_{x_0}.$$

The holonomy group of  $A$  at  $x_0$

$$\text{hol}_A(x_0) = \{\text{hol}_A(\gamma) \mid \gamma \in \Omega_{x_0}(X)\}$$

$\text{Stab}_A$  is a subgroup of  $G$  and is the commutant of the holonomy group of  $A$ . Note that the center of  $G$ ,  $Z(G) \subset \text{Stab}_A$  for all  $A$ .

### Definition

$A$  is called **reducible** if  $\text{Stab}_A \neq Z(G)$ .

## Exercise

*Show that the possible subgroups of  $SO_3$  that appear are  $1, \mathbb{Z}_2, \mathbb{Z}_2 \times \mathbb{Z}_2, SO_2, O_2$  and  $SO_3$  while if we restrict to the determinant 1 gauge group, the possible subgroups of  $SU_2$  that appear as  $\text{Stab}(A)$  for some connection are  $\mathbb{Z}_2, U_1$  and  $SU_2$ .*

The quotient space  $\mathcal{B} = \mathcal{A}/\mathcal{G}^1$  will have singularities if  $\mathcal{A}$  contains reducible connections. Compare

$$\mathfrak{so}_3^n / SO_3$$

where  $SO_3$  acts by the adjoint action. For each point in the quotient a neighborhood is modeled on a neighborhood of zero in one of the following three spaces.

We would like to show that  $\mathcal{B}$  has similar local models. ▶

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- $\mathfrak{so}_3^{n-1}$ .
- $(\mathbb{R}^2)^{n-1} / SO_2 \times \mathbb{R}^n$ .

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- $\mathfrak{so}_3^{n-1}$ .
- $(\mathbb{R}^2)^{n-1} / SO_2 \times \mathbb{R}^n$ .
- The origin which has no better model than  $\mathfrak{so}_3^n / SO_3$ .

We would like to show that  $\mathcal{B}$  has similar local models. ▶

We will need to work with Sobolev spaces and Fredholm operators on them. There are the very basic facts. Let  $L_k^p(T^n)$  be the completion of space of  $C^\infty$  functions with respect to the norm.

$$\|f\|_{L_k^p(T^n)}^p = \sum_{j=0}^k \int_{T^n} |\nabla^{(j)} f|^p dx^1 \wedge \dots \wedge dx^n.$$

We also will use Sobolev norms with fractional derivatives. These can be defined by interpolation. The  $L^2$ -version can be easily defined by Fourier transform. If

$$f(\mathbf{x}) = \sum_{\mathbf{n}} \hat{f}(\mathbf{n}) e^{i\mathbf{b} \cdot \mathbf{x}}$$

Then we

$$\|f\|_{L_s^2}^2 = \sum_{\mathbf{n}} (1 + |\mathbf{n}|^2)^{s/2} |\hat{f}(\mathbf{n})|^2$$

Then we have the following properties.

$$L_k^p(T^n) \hookrightarrow L_l^q(T^n)$$

If  $k - n/p \geq l - n/q$  and  $k \geq l$ . If both inequalities are strict the embedding is compact.

$$L_k^p(T^n) \hookrightarrow C^{l,\alpha}(T^n)$$

if  $k - n/p \geq l + \alpha$  where  $k \geq l \geq 0$  and  $\alpha > 0$ . If both inequalities are strict the embedding is compact.

Furthermore when  $k - n/p$  and  $l - n/q$  are both negative

$$L_k^p(T^n) \times L_l^q(T^n) \hookrightarrow L_m^r(T^n)$$

is continuous if

$$k - n/p + l - n/q \geq r - n/m \text{ and } k, l \geq r.$$

Suppose that  $k - n/p$  is positive and  $k, l \geq r$  then  $L_l^q(T^n)$  is module over  $L_k^p(T^n)$ .

In particular for  $n = 4$  we have

$$L_1^2 \hookrightarrow L^4,$$

$$L_2^2 \hookrightarrow L_1^4 \hookrightarrow L^p \text{ for all } p < \infty$$

$$L_2^2 \not\hookrightarrow C^0.$$

The following multiplications are continuous.

$$L_1^2 \times L_1^2 \mapsto L^2,$$

$$L_2^2 \times L_1^2 \mapsto L_1^p \text{ or } L_s^2 \text{ where } p < 2 \text{ or } s < 1.$$

Note that  $\nabla \ln(\ln(r)) = \frac{1}{r \ln(r)} = \frac{1}{r \ln(r)}$

$$\int_0^{1/2} \left(\frac{1}{r \ln(r)}\right)^4 r^3 dr = \int_0^{1/2} \left(\frac{1}{\ln(r)^4}\right) \frac{dr}{r} = -\frac{1}{4} \ln^{-3}(1/2) < \infty.$$

We need to complete our spaces of connections into Sobolev spaces. Suppose  $G = U_m$  or  $SO_m$  and let  $E$  be a  $\mathbb{C}^m$  or  $\mathbb{R}^m$  bundle so that  $P$  is the frame bundle of  $E$ . Using this connection there is a preferred gauge-equivariant Sobolev norm

$$\|f\|_{L_{k,A}^p}^p = \int_X \sum_{i=0}^k |\nabla_A^i f|^p *_g 1.$$

Any easy consequence of the Sobolev multiplication theorems is that the  $L_{k,A}^p$ -norm and the  $L_{k,A'}$  are equivalent norms provided that difference is in

$$A - A' \in L_{l,A}^q(X, T^*X \otimes \text{ad}P)$$

and the multiplication  $L_k^p \times L_l^q \rightarrow L_{k-1}^p$  is continuous, i.e.  $l \geq n/q - 1$  and  $l > k - 1$  and indeed there is a constant  $C > 0$  depending on  $\|A - A'\|_{L_{l,A}^q}$  so that

$$\|e\|_{L_{k,A'}^p} \leq C \|e\|_{L_{k,A}^p} \quad (1)$$

Fractional order Sobolev norms. Fix a smooth connection  $A_0$  and consider the corresponding Laplace operator

$$\nabla_{A_0}^* \nabla_{A_0} + 1 : C^\infty(X; \text{ad}P) \rightarrow C^\infty(X; \text{ad}P).$$

This operator has compact inverse and one can construct compact operators

$$(\nabla_{A_0}^* \nabla_{A_0} + 1)^{-s} : L^p(X; \text{ad}P) \rightarrow L^p(X; \text{ad}P)$$

for any  $s > 0$ . For smooth  $A_0$  these operators are represented by a kernel with pole of order  $\text{dist}(x, y)^{-n+2s}$  along the diagonal. The space of  $L_s^p$ -sections of a bundle is

$$\{(\nabla_{A_0}^* \nabla_{A_0} + 1)^{-s/2} \mathbf{e} \mid \mathbf{e} \in L^p(X; \text{ad}P)\}.$$

This space has a norm using the connection  $A_0$ ,

$$\|(\nabla_{A_0}^* \nabla_{A_0} + 1)^{-s/2} \mathbf{e}\|_{L_{s,A_0}^p} = \|\mathbf{e}\|_{L^p} \quad (2)$$



For any pair of connections  $A, A'$  with  $A' - A_0, A - A_0 \in L_{s'}^2$  where  $s' \geq n/q - 1$  and  $s' > s - 1$  the above definition of the spaces  $L_{s',A}^2$  and  $L_{s',A'}^2$  and the corresponding norms  $\|\cdot\|_{L_{s',A}^2}$  and  $\|\cdot\|_{L_{s',A'}^2}$  still make sense. Using that the fractional order spaces are interpolation spaces for the integral norms we see that the estimate (??) holds. When dependence of the norms on the choice of connection is not important for the discussion we will drop the connection from the notation for the norm. Note that if  $X$  is a manifold with boundary we can also define  $\dot{L}_{s,A}^p$  to be completion of smooth sections with support in the interior of  $X$  in the norm  $\|\cdot\|_{L_{s,A}^p}$ . Then we have the following useful duality result for  $1 \leq p < \infty$

$$(\dot{L}_{s,A}^p)^* = L_{-s,A}^q.$$

For consider the space  $\mathcal{A}_k^p$  of  $L_k^p$  connections. The gauge group that acts naturally on this space of connections is

$$\mathcal{G}_{k+1}^p = \{g \in L_{k+1}^p(X, \text{ad}P) \mid g \in \text{Ad}P \text{ a.e.}\}.$$

When  $k + 1 - n/p > 0$  this consists of **continuous** sections.

### Lemma

*For  $G = U_m$  or  $SO_m$  and  $k + 1 - n/p > 0$  is a Banach Lie group*

Proof. Consider the map

$$m : L_{k+1}^p(X, \text{End}(E)) \mapsto L_{k+1}^p(X, \text{End}(E))$$

given by

$$m(A) = AA^*$$

Then  $\mathcal{G}_{k+1}^p = m^{-1}(\mathbf{1})$ . For  $k + 1 - n/p > 0$   $m$  is a smooth map since  $L_{k+1}^p$  is a Banach algebra.  $M$  is also easily seen to have surjective differential and the kernel of

## Lemma

For  $k + 1 - n/p > 0$ , The map  $\mathcal{A}_k^p \times \mathcal{G}_{k+1}^p \rightarrow \mathcal{A}_k^p$  is a smooth map of Banach manifolds.

Proof: In a trivialization the map is given by

$$(a, g) \mapsto gag^{-1} + gdg^{-1}$$

By the assumptions  $L_{k+1}^p$  is a Banach algebra and  $L_k^p$  is module over  $L_{k+1}^p$ . The above map is a composition of a continuous linear maps and continuous multiplications so it is smooth.  $\square$

The space of connections, being an affine space is always a Banach manifold.

### Lemma

For  $k + 1 - n/p \geq 0$ ,  $\mathcal{B}_k^p = \mathcal{A}_k^p / \mathcal{G}_{k+1}^p$  is a Hausdorff topological space.

Proof: We prove this in the case that of  $L_2^2$  gauge transformations. The general case is similar. We must show that  $\{(A, g \cdot A) \mid A \in \mathcal{A}_1^2, g \in \mathcal{G}_2^2\} \subset \mathcal{A}_1^2 \times \mathcal{A}_1^2$  is closed. Fix  $A_0 \in \mathcal{A}_k^p$

$$(A_i, g_i \cdot A_i) \xrightarrow{i \rightarrow \infty} (A, B)$$

$$A_i = A_0 + a_i, g_i \cdot A_i = A_0 + b_i.$$

Then

$$g_i d_{A_0} g_i^{-1} + g_i a_i g_i^{-1} = b_i \Rightarrow$$

$$d_{A_0} g_i = g_i a_i - b_i g_i.$$

(3)



The sequences  $a_i, b_i$  converge in  $L^2_1$  to  $a$  and  $b$  respectively. We need to show that  $A_0 + a$  is  $L^2_2$  gauge equivalent to  $A_0 + b$ . We see that  $d_{A_0} g_i$  bounded in  $L^4$  so  $g_i$  is bounded in  $L^4_1$ . Now differentiate Equation (??)

$$\nabla_{A_0}^2 g_i = \nabla_{A_0} g_i a_i - b_i \nabla_{A_0} g_i + g_i \nabla_{A_0} a_i - \nabla_{A_0} b_i g_i$$

The right hand side is bounded in  $L^2$  for each is either  $L^2 \times L^\infty$  or  $L^4 \times L^4$ . Thus we can pass to subsequence where  $g_i$  converges  $L^2_2$ -weakly to  $g$  and strongly in  $L^2$ . The strong convergence implies again after passing to a subsequence a.e. convergence and so  $g g^* = 1$  a.e. thus  $g$  is gauge transformation. We can take the limit in Equation (??) in  $L^r$  for  $r < 2$  showing that  $g$  takes  $A + a$  to  $A + b$  as required.  $\square$ .

**Orbifold structure for  $\mathcal{A}/\mathcal{G}$ .** For concreteness consider we'll restrict to the  $L^2$  case and consider the space,  $\mathcal{A}_{3/2}$ , of connections on a four-manifold  $X$ . The gauge group,  $\mathcal{G}_{5/2}$ . Slices. Take orthogonal complement of tangent space of gauge orbit.

$$T_A\mathcal{G} = d_A\Omega^0(X, \text{ad}P).$$

Using  $d \text{tr}(\sigma \wedge *a) = \text{tr}(d_A\sigma \wedge *a) + \text{tr}(\sigma \wedge d_A *a)$  we have

$$-\int_X \text{tr}(d_A\sigma \wedge *a) = \int_X \text{tr}(\sigma \wedge d_A *a) = -\int_X \text{tr}(\sigma \wedge *d_A^*a).$$

Thus we want

$$0 = \langle d_A\sigma, a \rangle = \langle \sigma, d_A^*a \rangle.$$

for all  $\sigma \in \Omega^0(X; \text{ad}P)$  so  $d_A^*a = 0$ .

$A \in \mathcal{A}_{3/2}$  we have:

$$\mathcal{S}_{S,A}(\epsilon) = \{A + a \mid a \in L^2_{3/2,A}(X; T^*X \otimes \text{ad}P),$$

$$d_A^*a = 0 \text{ and } \|a\|_{L^2_{S,A}} < \epsilon\}$$

To state the result in general we need to discuss how to deal with reducible connections. If  $A$  is reducible  $\text{Stab}_A$ , the stabilizer of  $A$ , is a Lie subgroup of  $SO_m$ .  $\text{Stab}_A$  acts on  $S_A$  and freely on  $\mathcal{G}_{5/2}$  and it is straightforward to see that the quotient

$$(\mathcal{S}_{A,3/2}(\epsilon) \times \mathcal{G}_{5/2})/\text{Stab}_A$$

is a smooth Hilbert manifold. The tangent space at the equivalence class of  $(A + a, 1)$  is identified with the quotient of

$$\{\alpha \mid \alpha \in L_{3/2,A}^2, d_A^* \alpha = 0\} \times L_{5/2,A}^2(X, \text{ad}P)$$

by the finite-dimensional (and hence closed) subspace

$$\{(-[a, \xi^0], \xi^0) \mid d_A \xi^0 = 0\}.$$

## Proposition

*For all  $A \in \mathcal{A}_{3/2}$  there is an  $\epsilon > 0$  so that the map*

$$\begin{aligned} m &: (\mathcal{S}_{3/2,A}(\epsilon) \times \mathcal{G}_{5/2}) / \text{Stab}_A \rightarrow \mathcal{A}_{3/2} \\ m([A + a, g]) &= A + gag^{-1} - (d_A g)g^{-1} \end{aligned}$$

*is a  $\mathcal{G}_{5/2}$ -equivariant diffeomorphism onto its image.*



The proof will use the following facts.

- The operator

$$d_A : L_{A,s}^p(X, \text{ad}P) \rightarrow L_{A,s-1}^p(X, T^*X \otimes \text{End}(E))$$

has finite dimensional kernel and closed range.

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- There is a closed "Hodge"  $L^2$  orthogonal decomposition.

$$L_{A,s}^p(X, \text{ad}P) = \ker(d_A) \oplus d_A^*(L_{A,s+1}^p(X, T^*X \otimes \text{ad}P))$$

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- A complement for the range of

$$d_A : L_{A,s}^p(X, \text{ad}P) \rightarrow L_{A,s-1}^p(X, T^*X \otimes \text{End}(E)) \text{ is}$$

$$\ker d_A^* : L_{A,s-1}^p(X, T^*X \otimes \text{ad}P) \rightarrow L_{A,s-2}^p(X, \text{ad}P)$$

The proof will use the following facts.

- The operator

$$d_A : L_{A,s}^p(X, \text{ad}P) \rightarrow L_{A,s-1}^p(X, T^*X \otimes \text{End}(E))$$

has finite dimensional kernel and closed range.

- There is a closed "Hodge"  $L^2$  orthogonal decomposition.

$$L_{A,s}^p(X, \text{ad}P) = \ker(d_A) \oplus d_A^*(L_{A,s+1}^p(X, T^*X \otimes \text{ad}P))$$

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- There is a constant  $\lambda_1(A) > 0$  so that for all  $\xi \perp \ker(d_A)$  we have  $\int_X |d_A \xi|^2 \geq \lambda_1 \int_X |\xi|^2$ .

Proof.

Write  $\xi = \xi^0 + \xi^1$  for the Hodge decomposition where  $d_A \xi^0 = 0$  and  $\xi^1$  is  $L^2$ -orthogonal to elements of the kernel of  $d_A$ .

$m$  is a local diffeomorphism if the differential of  $m$  is an isomorphism at  $(A, 1)$ . Since  $m$  is  $\mathcal{G}_{5/2}$ -equivariant it suffices to check this at the equivalence classes of  $(A, 1)$ . The differential is given by the map

$$\mathcal{D}_{(A,1)}m([\alpha, \xi]) = d_A \xi + \alpha$$

This map is surjective by the Hodge decomposition. If  $(\alpha, \xi)$  is in the kernel then  $\alpha = 0$  and  $\xi \in \ker d_A$  so  $[\alpha, \xi] = 0$ .

$m$  is injective. It suffices to show that if  $g \cdot (A + a) = A + b$  is in the slice then  $g \in \text{Stab}_A$ . The condition that  $g \cdot (A + a) = A + b$  is equivalent to

$$d_A g = ga - bg. \quad (4)$$

Taking  $d_A^*$  of this equation gives, using the slice condition  $d_A^* a = d_A^* b = 0$ ,

$$d_A^* d_A g = - * d_A g \wedge * a - * b \wedge * d_A g. \quad (5)$$

Use the Hodge decomposition  $g = g^0 + g^1$  again and take the inner product of this equation with  $g^1$ .

$$\begin{aligned}
\|d_A g^1\|_{L^2}^2 &\leq (\|a\|_{L^4} + \|b\|_{L^4}) \|d_A g^1\|_{L^2} \|g^1\|_{L^4} \\
&\leq \kappa (\|a\|_{L^4} + \|b\|_{L^4}) \|d_A g^1\|_{L^2} \|g^1\|_{L^2_{1,A}} \\
&\leq \kappa (1 + \lambda_1^{-1}(A))^{1/2} (\|a\|_{L^4} + \|b\|_{L^4}) \|d_A g^1\|_{L^2}^2
\end{aligned}$$

We have used

$$\begin{aligned}
\|g^1\|_{L^2_{1,A}}^2 &= \|g^1\|_{L^2}^2 + \|d_A g^1\|_{L^2}^2 \\
&\leq (\lambda_1(A)^{-1} + 1) \|d_A g^1\|_{L^2}^2.
\end{aligned}$$

So for  $\epsilon \kappa (1 + \lambda_1^{-1}(A))^{1/2} < 1/2$  we have  $u^1 = 0$  and hence  $u = u^0$  is in  $\text{Stab}(A)$  as required.  $\square$

## Remark

*Not that the proof of injectivity of  $\mathfrak{m}$  only used  $L^4$ -smallness. **exercise** Repeat this argument for  $\mathcal{D}_{(1,A)}\mathfrak{m}$  so show that for all  $A \in \mathcal{S}_{1,A}$  this differential is invertible. This then can be used to give a sharper result that will be important in the proof of Uhlenbeck compactness.*



## Remark

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## Proposition (Big slices)

For all  $A \in \mathcal{A}_{3/2}$  then the map

$$\begin{aligned} m &: (\mathcal{S}_{1,A}(\epsilon) \times \mathcal{G}_{5/2}/\text{Stab}_A) \rightarrow \mathcal{A}_{3/2} \\ m([A + a, g]) &= A + gag^{-1} - (d_A g)g^{-1} \end{aligned}$$

is a  $\mathcal{G}_{5/2}$ -equivariant diffeomorphism onto its image. The image contains an  $L^2_1$  ball about  $A$ .

When we have a manifold with boundary the notion of slice needs to be modified.

$$\int_X \text{tr}(d_A \sigma \wedge *a) = \int_{\partial X} \text{tr}(\sigma \wedge *a) + \int_X \text{tr}(\sigma \wedge *d_A^* a).$$

So the condition of being orthogonal to the tangent space of the gauge group implies formally

$$d_A^* a = 0 \text{ and } *a|_{\partial X} = 0.$$

In this and the next lecture we prove Uhlenbeck's Compactness Theorem for Yang-Mills connections.

### Lemma

Let  $B$  be a geodesic ball in a Riemannian manifold. Let  $\Gamma$  be the connection arising from a trivialization  $P = \times G$  for principal  $G$  bundle. There are positive constants  $C, \epsilon_0$  so that for all  $A = \Gamma + a$  with  $a \in L^2_1(B, T^*B \times \mathfrak{g})$  and  $\int_B \text{tr}(F_A \wedge *F_A) \leq \epsilon_0$  there is a map  $g \in L^2_2(B, M_n(\mathbb{C}))$  with  $g \in G$  a.e. so that if we set

$$b = gag^{-1} + gdg^{-1}$$

we have

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$$b = gag^{-1} + gdg^{-1}$$

we have

- $d^*b = 0$  and  $*b|_{\partial B} = 0$  and
- $\int_B (|\nabla_\Gamma b|^2 + |b|^2) * 1 \leq C \int_B |F_A|^2 * 1.$

Proof.  $B = \exp_x(B_0(\delta))$ . The first move in the proof is notice that we can replace the  $L^2$ -norm of the curvature by an equivalent norm. We exploit this possibility by changing the metric to the pullback metric from the metric from  $T_x B$ . Let  $\tilde{*}$  denote the Hodge star for this metric. The proof of this is a version of the continuity argument. Let

$$V_\epsilon = \{a \in L^2_{\frac{3}{2}}(B, T^*B \otimes \mathfrak{g}) \mid - \int_B \text{tr}(F_A \wedge \tilde{*}F_A) \leq \epsilon\}.$$

We use a density argument to get the  $L^2_1$ -case. Let  $W_\epsilon \subset V_\epsilon$  be the subset where there is a  $C > 0$  so that conclusions of the lemma holds.

We will show that for  $\epsilon$  small enough and  $C$  large enough that:

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Thus  $W_\epsilon = V_\epsilon$ .

$V_\epsilon$  is connected. Let for  $0 \leq t \leq 1$  let  $\tau_t : B_x(\delta) \rightarrow B_x(t\delta)$  denote the scaling by  $t$  along geodesics through  $x$ .

Consider the path of connection

$$A_t = \tau_t^*(A|_{B_x(t\delta)}).$$

We claim that for  $A$  an  $L^2_{3/2}$ -connection this path is continuous in the  $L^2_{3/2}$ -topology. For  $t \neq 0$  this is the same point that translation is point-wise continuous on  $L^p$  spaces. At  $t = 0$  continuity follows directly from the definitions.

Since the  $L^2$ -norm of the curvature is conformally invariant and dilation is a conformal map we have that

$$-\int_B (\text{tr}(F_{A_t} \wedge \tilde{*}F_{A_t})) = -\int_{B_x(t\delta)} (\text{tr}(F_A \wedge \tilde{*}F_A)) \leq \epsilon$$

and  $\lim_{t \rightarrow 0} -\int_B (\text{tr}(F_{A_t} \wedge \tilde{*}F_{A_t})) = 0$ .

$W_\epsilon$  is closed in  $V_\epsilon$ . Let  $A_i = \Gamma + a_i$  be a sequence in  $W_\epsilon$  which converges in the  $L^2_{3/2}$ -topology to  $A = \Gamma + a \in V_\epsilon$ . There is also then a sequence of gauge transformations  $g_i$  so that  $g_i \cdot A_i = \Gamma + b_i$  and the  $b_i$  are in Coulomb gauge and satisfy the estimates, so it follows that  $b_i$  converge weakly in the  $L^2_1$  topology to  $b$ . Now consider the sequence of gauge transformations  $g_i$ . They satisfy the equations:

$$dg_i = g_i a_i - b_i g_i.$$

Thus  $dg_i \in L^4$  and so  $g_i \in L^4_1$  ( $g_i$  is bounded). Taking the gradient of both sides gives:

$$\nabla dg_i = \nabla g_i a_i + g_i \nabla a_i - \nabla b_i g_i - b_i \nabla g_i$$

hence  $g_i$  is bounded in  $L^2_2$ . Thus after passing to a subsequence we may assume the  $g_i$  converge weakly in  $L^2_2$  to  $g$ .

$$|g_i| \leq M \xrightarrow{\text{subseq}} g_i \xrightarrow{\text{a.e.}} g.$$

Hence  $g \in U_n$  almost everywhere.

$$\|a_i\|_{L^2_1}, \|b_i\|_{L^2_1}, \|g_i\|_{L^2_2} \leq M \xrightarrow{\text{subseq}} a_i \xrightarrow{L^4} a, b_i \xrightarrow{L^4} b, g_i \xrightarrow{L^4, L^2_1} g.$$

where  $a_i, b_i$  converge strongly in  $L^{4-\epsilon}$  and  $g_i$  converges strongly in  $L^p$  for all  $p$  and in  $L^2_1$  so that in the equation

$$dg_i = g_i a_i - b_i g_i.$$

both sides converges strongly in  $L^2$ . Thus the we have

$$dg = ga - bg$$

so that  $\Gamma + a$  is gauge equivalent to  $\Gamma + b$ . To see that  $\Gamma + a \in W_\epsilon$  we need to see that  $b$  satisfies the estimate and that  $b \in L^2_{3/2}$

$$\begin{aligned}
\int_U |F_{\Gamma+b}|^2 &= \int_U \langle db + b \wedge b, db + b \wedge b \rangle \\
&\geq \int_U |db + d^*b|^2 - \|b\|_{L^4}^4 \\
&\geq \int_U |\nabla b|^2 - C\epsilon \|b\|_{L^2_1}^2 \\
&\geq ((1 + \lambda_1)/2 - C\epsilon) \|b\|_{L^2_1}^2.
\end{aligned}$$

We've used the identity

$$\int_U (|(d + d^*)b|^2) = \int_U |\nabla b|^2 + \int_{\partial U} *bd$$

Check that  $b \in L^2_{3/2}$ . Take  $d^*$  of this equation giving

$$\Delta g = - * (dg \wedge *a) + gd^*a + *(*b \wedge dg).$$

Rewrite as:

$$\Delta g - *(*b \wedge dg) = - * (dg \wedge *a) + gd^*a.$$

Note that we also have  $*dg|_{\partial B} = *a|_{\partial b}$ . The red terms both are in  $L^2_{1/2}$  since  $L^2_{3/2} \times L^2_1 \mapsto L^2_{1/2}$ . Thinking of the blue term as function of  $g$  it is a bounded operator

$$L^2_{5/2} \rightarrow L^2_{1/2}$$

with operator norm controlled by  $\|b\|_{L^2_1}$ . The map  $h \mapsto (\Delta h, *dh|_{\partial U})$  viewed a map

$$L^2_{5/2}(U) \rightarrow L^2_{1/2}(U) \times L^2_1(\partial U)$$

is surjective with finite dimensional kernel.



The equation has index 0 and is invertible for  $b = 0$  so for  $b$  with small  $L_1^2$ -norm it follows that  $g \in L_{5/2}^2$  is the unique solution.

The equation

$$b = gag^{-1} - g^{-1}dg$$

says that  $b$  is in  $L_{3/2}^2$ .

$W_\epsilon$  is open in  $V_\epsilon$  Openness follows from the Big Slices Lemma, Lemma ?? which tells us that if  $A$  is gauge equivalent into the  $S_{r,1}(\epsilon)$  then so is a neighborhood of  $A$  in the  $L^2_{3/2}$ -topology.

Finally suppose that  $A_i$  is a sequence of  $L^2_{3/2}$  connections in  $V_\epsilon$  converging to an  $L^2_1$  connection  $A$  then following the proof of closedness we see that  $A$  can be put into a good gauge and this representative satisfies the estimates.  $\square$

The same circle of ideas proves the following lemma which will later be used in the proof of the Uhlenbeck compactness theorem.

### Lemma (Curvature is Proper)

With  $\epsilon_0$  as in the previous lemma. Let  $SV_{(\epsilon_0/2)}$  be the set of  $L^2_1$  connections  $A = \Gamma + a$  with

$$\int_B |F_A|^2 * 1 \leq \epsilon_0/2$$

and

$$d^* a = 0 \text{ and } * a|_{\partial B} = 0.$$

Then the curvature map

$$\begin{aligned} F_\bullet : SV_{\epsilon_0/2} &\rightarrow L^2(B, \Lambda^2(T^*B) \otimes \text{ad}P) \\ A &\rightarrow F_A \end{aligned}$$

is proper

Proof. Suppose  $F_{A_i}$  is a Cauchy sequence.

$$\begin{aligned} \int_B |F_{A_i} - F_{A_j}|^2 &= \int_B |da_i + a_i \wedge a_i - da_j - a_j \wedge a_j|^2 * 1 \\ &\geq \int_B (|da_i - da_j|^2 - |a_i \wedge a_i - a_j \wedge a_j|^2) * 1 \\ &\geq \int_B |(d + d^*)(a_i - a_j)|^2 - (\|a_i\|_{L^4} + \|a_j\|_{L^4}) \|a_i - a_j\|_{L^4}^2 \\ &\geq \int_B |\nabla(a_i - a_j)|^2 - 2C\epsilon \|a_i - a_j\|_{L^2_1}^2 \\ &\geq ((1 + \lambda)/2 - 2C\epsilon) \|a_i - a_j\|_{L^2_1}^2. \end{aligned}$$

Thus  $A_i$  is also Cauchy.  $\square$

$X$  an oriented riemannian  $n$ -manifold. The **Hodge Star**

$$* : \Lambda^k \rightarrow \Lambda^{n-k}$$

Behavior under conformal change:

$$g \mapsto e^{2\sigma} g \implies * \mapsto e^{(n-2k)\sigma} *$$

In particular if  $n = 2k$  Hodge star is **conformally invariant**.  
Also note that

$$*^2 = (-1)^{k(n-k)}$$

In particular for  $n = 4, k = 2$

$$* : \Lambda^2 \rightarrow \Lambda^2$$

has  $*^2 = 1$ .

Thus 2-forms have a conformally invariant decomposition.

$$\begin{aligned}\Lambda^2 &= \Lambda^+ \oplus \Lambda^- \\ * &= 1 \quad * = -1\end{aligned}$$

Explicitly:  $e^1, e^2, e^3, e^4$  oriented orthonormal frame.

$$\omega_I = e^1 \wedge e^2 + e^3 \wedge e^4, \omega_J = e^1 \wedge e^3 - e^2 \wedge e^4, \omega_K = e^1 \wedge e^4 + e^2 \wedge e^3$$

span  $\Lambda^+$ , the **Self-Dual** two forms and

$$e^1 \wedge e^2 - e^3 \wedge e^4, e^1 \wedge e^3 + e^2 \wedge e^4, e^1 \wedge e^4 - e^2 \wedge e^3$$

span  $\Lambda^-$ , the **Anti-Self-Dual** two forms. Note that  $\Lambda^+$  and  $\Lambda^-$  are pointwise orthogonal under **both** the Riemannian inner product and the wedge product.

Thus we can decompose the curvature of a connection:

$$F_A = F_A^+ + F_A^-.$$

$F_A^+ = 0$  ( $F_A^- = 0$ ) are the Anti-Self-Dual (Self-Dual) Yang-Mills equation.

Note that the Bianchi identity implies

$$d_A F_A = 0$$

and hence if  $A$  is SD or ASD we have

$$d_A^* F_A = - * d_A * F_A = \pm * d_A F_A = 0.$$

Thus SD and ASD are critical points for  $E$  but more is true!

If  $X$  is a closed four-manifold recall:

$$8\pi^2 k = \int_X \text{tr}(F_A \wedge F_A), \quad E(A) = - \int_X \text{tr}(F_A \wedge *F_A) \geq 0.$$

Thus adding and subtracting these formulae gives:

$$E(A) + 8\pi^2 k = -2 \int_X \text{tr}(F_A^- \wedge *F_A^-) \geq 0$$

and

$$E(A) - 8\pi^2 k = -2 \int_X \text{tr}(F_A^+ \wedge *F_A^+) \geq 0$$

Thus

$$E(A) \geq 8\pi^2 |k|$$

with equality if and only if  $F_A^+ = 0$  when  $k \geq 0$  and  $F_A^- = 0$  when  $k \leq 0$ .



The basic instanton. We'll use quaternionic notation.  
 $\mathbb{H} = \mathbb{R} + \mathbb{R}I + \mathbb{R}J + \mathbb{R}K$  where  $IJ = K + \text{cyclic}$  and  
 $IJ = -JI + \text{cyclic}$ .

$$\mathbf{x} = x_0 + x_1I + x_2J + x_3K$$

Conjugation

$$\bar{\mathbf{x}} = x_0 - x_1I - x_2J - x_3K.$$

$S^7 \subset \mathbb{H}^2$  the unit sphere

Two different  $SU_2 = Sp(1) = \{\mathbf{x} \in \mathbb{H} | \mathbf{x}\bar{\mathbf{x}} = \mathbf{1}\}$  actions.

$$(\mathbf{x}_0, \mathbf{x}_1)q = (\mathbf{x}_0q, \mathbf{x}_1q) \text{ or}$$

$$(\mathbf{x}_0, \mathbf{x}_1)q = (\bar{q}\mathbf{x}_0, \bar{q}\mathbf{x}_1).$$

In either case the quotient is  $S^4 = \mathbb{H}P^1$ .



Two principal bundle with total space  $S^7$ .

$$P_{\pm} \rightarrow S^4$$

where  $P_+$  has the right action and  $P_-$  has the left action.  
These are the unit sphere bundles of  $\mathbb{H}$ -bundles  $S_{\pm} \rightarrow S^4$ .  
For example

$$S_+ = \{([\mathbf{x}_0, \mathbf{x}_1], (\mathbf{v}_0, \mathbf{v}_1)) \mid (\mathbf{v}_0, \mathbf{v}_1) = \bar{h}(\mathbf{x}_0, \mathbf{x}_1), \quad h \in \mathbb{H}\}.$$

Give  $P_{\pm}$  the connection  $A_{\pm}$  which declares that the horizontal space at  $p \in P_{\pm}$  is the orthogonal complement of the fiber  $P_{\pm} \rightarrow S^4$ .  $Sp_2(= Spin_5)$  acts on  $S^7$  isometrically preserving the connections  $A_{\pm}$ . The stabilizer of a point  $p \in S^7$  is a copy of  $Sp_1(= SU_2 = Spin_3)$  so that the curvature of the connection is an  $SU_2$  equivariant map

$$\Lambda^2(T_p^*S^4) = \Lambda^+ \oplus \Lambda^-(T_p^*S^4) \rightarrow \text{ad}P_{\pm}|_p.$$

Need to understand  $SU_2$  action.  $\text{ad}P|_{\pm}$  is the vertical tangent space at  $p$  which is canonically identified with  $\mathfrak{su}_2$  with the adjoint action.

Write  $d\mathbf{x} = dx_0 + dx_1 I + dx_2 J + dx_3 K$ . Note that  $d\mathbf{x} \wedge d\bar{\mathbf{x}}$  is a purely imaginary 2-form (*im*  $\mathbb{H}$ -valued). Indeed

$$d\mathbf{x} \wedge d\bar{\mathbf{x}} = -2(\omega_I I + \omega_J J + \omega_K K).$$

In particular  $d\mathbf{x} \wedge d\bar{\mathbf{x}}$  is self-dual.

$$d\bar{g}\mathbf{x}h \wedge d\overline{\bar{g}\mathbf{x}h} = \bar{g}d\mathbf{x} \wedge d\bar{\mathbf{x}}g.$$

Thus  $d\mathbf{x} \wedge d\bar{\mathbf{x}}$  is invariant under the left action and equivariant under the right action). Thus  $A_+$  is a self-dual connection.

Similarly  $d\bar{\mathbf{x}} \wedge d\mathbf{x}$  is anti-self-dual and invariant under the left action and equivariant under the right action. Thus  $A_-$  is an anti-self-dual connection.

Exercise: Show that the connection one-form

$$\begin{aligned} a &= \frac{\operatorname{Im} \bar{\mathbf{x}} \wedge d\mathbf{x}}{(1 + |\mathbf{x}|^2)} \\ &= \frac{\bar{\mathbf{x}} \wedge d\mathbf{x} - \mathbf{x} \wedge d\bar{\mathbf{x}}}{2(1 + |\mathbf{x}|^2)} \end{aligned}$$

represents  $A_+$  in a suitable trivialization of  $P_+$ .

We can construct from  $A_-$  other ASD-connections using conformal invariance. The dilatation  $\tau_\lambda(\mathbf{x}) = \lambda\mathbf{x}$  induces a conformal diffeomorphism. Indeed the basic instanton is invariant under  $SO_5$  acting on  $S^4$  but the conformal group  $SO_{5,1}$  acts so effectively there is a

$$SO_{5,1}/SO_5 = \mathbb{H}^5$$

worth of ASD-connections in  $P_+$ . Atiyah-Hitchin-Singer prove any ASD connections in  $P_+$  is gauge equivalent to one of these.

This example exhibits the phenomenon of **bubbling**. As the the parameter  $\lambda \mapsto \infty$  the connections

$$\begin{aligned}\tau_\lambda^*(a) &= \frac{\operatorname{Im} \lambda \bar{\mathbf{x}} \wedge d\lambda \mathbf{x}}{(1 + |\lambda \mathbf{x}|^2)} \\ &= \frac{\operatorname{Im} \bar{\mathbf{x}} \wedge d\mathbf{x}}{(1/\lambda^2 + |\mathbf{x}|^2)}\end{aligned}$$

converge away from the origin to

$$\frac{\operatorname{Im} \bar{\mathbf{x}} \wedge d\mathbf{x}}{(|\mathbf{x}|^2)}$$

which is gauge equivalent to zero! by the gauge transformation

$$g(\mathbf{x}) = \mathbf{x}/|\mathbf{x}|.$$

An  $SO_3$  bundle on  $P \rightarrow X$  has two characteristic classes

$$p_1(P) \in H^4(X; \mathbb{Z}) \cong \mathbb{Z} \text{ and } w_2(P) \in H^2(X; \mathbb{Z}_2)$$

These are constrained by

$$\wp(w_2(P)) \equiv p_1(P) \pmod{4}.$$

where

$$\wp : H^2(X; \mathbb{Z}_2) \rightarrow H^4(X; \mathbb{Z}_4)$$

is the Pontryagin square.

Given  $P \rightarrow X$  and riemannian metric  $g$  let

$$M_{k,w} = \{A \in \mathcal{A}(P) \mid F_A^+ = 0\} / \mathcal{G}(P).$$

where  $k = -\frac{1}{4} \langle p_1(P), [X] \rangle$  and  $w = w_2(P) \in H^2(X; \mathbb{Z}_2)$ .



## Fredholm maps and Kuranshi description of the moduli space.

### Definition

A smooth map of Hilbert manifolds  $f : \mathcal{M} \rightarrow \mathcal{N}$  is called Fredholm if for all  $x \in \mathcal{M}$

$$d_x f : T_x \mathcal{M} \rightarrow T_{f(x)} \mathcal{N}$$

is a Fredholm map.

## Theorem

*If  $f : \mathcal{M} \rightarrow \mathcal{N}$  is a Fredholm map and  $x \in \mathcal{M}$  then there is*

*So that  $\phi(0) = x$ ,  $d_0\phi : T_0U \rightarrow \ker(d_x f) \subset T_x\mathcal{M}$  and so that*

$$\phi(\kappa^{-1}(0)) = f^{-1}(0) \cap \mathcal{W}.$$

## Theorem

If  $f : \mathcal{M} \rightarrow \mathcal{N}$  is a Fredholm map and  $x \in \mathcal{M}$  then there is

- a neighborhood  $0 \ni U$   $\ker(d_x f) \subset T_x \mathcal{M}$  and neighborhood  $x \ni \mathcal{W} \subset \mathcal{M}$ .

So that  $\phi(0) = x$ ,  $d_0 \phi : T_0 U \rightarrow \ker(d_x f) \subset T_x \mathcal{M}$  and so that

$$\phi(\kappa^{-1}(0)) = f^{-1}(0) \cap \mathcal{W}.$$

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If  $A$  is an ASD connection then we have the **deformation complex**

$$0 \mapsto L_3^2(X; \text{ad}P) \xrightarrow{d_A} L_2^2(X; \Lambda^1 \otimes \text{ad}P) \xrightarrow{d_A^+} L_1^2(X; \Lambda^+ \otimes \text{ad}P) \rightarrow 0.$$

This is an elliptic complex and the index of the complex is

$$-8k + \frac{3}{2}(\chi(X) + \sigma(X))$$

We call the negative of the index

$$8k - \frac{3}{2}(\chi(X) + \sigma(X))$$

the **formal dimension** of the moduli space.

This index was computed by Atiyah-Hitchin-Singer using the index theorem. The result is.

## Theorem

*The map  $A + a \mapsto F^+(A)$  is a Fredholm map viewed as a map*

$$\phi : \mathcal{S}_{A,2} \rightarrow L_1^2(X, \Lambda^+(X)).$$

*The index of  $\phi$  is*

$$8k - \frac{3}{2}(\chi(X) + \sigma(X)) + \dim \text{Stab}(A)$$

## Examples.

1.  $X = S^4$ ,  $k$ ,  $M_{1,0}(S^4) = B^5$  and dimension formula gives

$$8k - \frac{3}{2}(2 + 0) = 8k - 3$$

Note that  $k = 1$  case checks with the example.

2.  $X$  arbitrary,  $k = 0$ .  $A$  is a trivial connection. The complex becomes 3 copies of

$$0 \mapsto L^2_3(X) \xrightarrow{d} L^2_2(X; \Lambda^1) \xrightarrow{d^+} L^2_2(X; \Lambda^+) \rightarrow 0.$$

where the dimensions of the cohomology groups are  $b^0(X)$ ,  $b^1(X)$ ,  $b^+(X)$  respectively.

$$\begin{aligned} & -\frac{3}{2}(\chi(X) + \sigma(X)) \\ &= -\frac{3}{2}(2b^0(X) - 2b^1(X) + b^+(X) + b^-(X) + b^+(X) - b^-(X)) \\ &= -3(b^0(X) - b^1(X) + b^+(X)). \end{aligned}$$



These calculations can be turned around to prove the dimension formula using *excision*.

**The Chern-Simons Function.** Given  $A$  connection in  $P \rightarrow X$ .

$$\int_X \text{tr}(F_A \wedge F_A) \pmod{8\pi^2\mathbb{Z}}$$

depends only on the **gauge equivalence class** of  $A|_{\partial X}$ .  
Indeed if  $A'$  is a connection in  $P'$  and there is a bundle iso  $g : P|_{\partial X} \rightarrow P'|_{\partial X}$  so that  $gA = A'$  then

$$\int_X \text{tr}(F_A \wedge F_A) - \int_X \text{tr}(F_{A'} \wedge F_{A'}) = \int_{DX} \text{tr}(F_{A''} \wedge F_{A''}) \in 8\pi^2\mathbb{Z}.$$

$DX$  is the double of  $X$  and  $A''$  is a connection in a bundle  $P'' \rightarrow DX$  obtained by patching together  $P$  with the connection  $A$  and  $P'$  with the connection  $A'$  using  $g$ .

## Definition

Let  $B$  be a connection in  $Q \rightarrow Y$

$$CS(B) = \int_X \text{tr}(F_A \wedge F_A) \pmod{8\pi^2\mathbb{Z}}$$

Another viewpoint on  $CS$ . Given a pair of connections  $B_0$  and  $B_1$  choose  $B_t$ ,  $t \in [0, 1]$ . View this path as a connection  $A$  in  $[0, 1] \times Q \rightarrow [0, 1] \times Y$ .

$$CS_{B_0}(B_1) = \int_{[0,1] \times Y} \text{tr}(F_A \wedge F_A).$$

If  $B_0$  comes from a trivialization and  $B_1 = B_0 + b$  and we set  $B_t = B_0 + tb$  so that:

$$F_A = d(tb) + \frac{t^2}{2}[b \wedge b] = dt \wedge b + tdb + \frac{t^2}{2}[b \wedge b].$$

Then

$$\text{tr}(F_A \wedge F_A) = dt \wedge \text{tr}(2b \wedge (tdb + t^2[b \wedge b]))$$

and so

$$\begin{aligned} \int_{[0,1] \times Y} \text{tr}(F_A \wedge F_A) &= \int_{[0,1] \times Y} dt \wedge \text{tr}(2b \wedge tdb + t^2 b \wedge [b \wedge b]) \\ &= \int_Y \text{tr}(b \wedge db + \frac{1}{3} b \wedge b \wedge b). \end{aligned}$$

## Exercise

*If  $B_0$  does not arise from a trivialization show that*

$$CS_{B_0}(B_1) = \int_Y \text{tr}(2b \wedge F_{B_0} + b \wedge db + \frac{1}{3} b \wedge [b \wedge b])$$

## Proposition

*Suppose  $G = SU_2$ . Then  $P = Y \times SU_2$  and  $g$  can be viewed as a map  $g : Y \rightarrow SU_2$ . We claim*

$$CS(B_1) - CS(gB_1) = \deg(g)$$

Proof: In this case the difference is the Chern-Weil integral for a connection in the bundle arising from the mapping torus of  $g$ . The  $c_2(P)$  is the Euler class of the associated complex two plane bundle. There is a section of this mapping torus bundle with  $\deg(g)$  transverse zeroes.

The Chern-Simons functions provides another local characterization of Anti-Self-Duality.

### Proposition (Mean Value Property)

Suppose  $M$  is connected smooth manifold with boundary.  $A$  in  $P \rightarrow X$  the

$$-\int_X \text{tr}(F_A \wedge *F_A) \geq -\text{CS}(A|_{\partial X})$$

with equality if and only if  $A$  is ASD.

Proof:

$$-\int_X \text{tr}(F_A \wedge *F_A) + \int_X \text{tr}(F_A \wedge F_A) = -2 \int \text{tr}(F_A^+ \wedge *F_A^+).$$



Uhlenbeck's compactness theorem. Let  $U$  be a small geodesic ball in a Riemannian four-manifold. Since all bundles on a ball are trivial we write  $M(U)$  for the moduli space on  $U$ . For  $\epsilon > 0$  write

$$M^\epsilon(U) = \{[A] \in M(U) \mid \int_U |F_A|^2 * 1 \leq \epsilon\}.$$

### Proposition

*Let  $U$  be a geodesic ball in a Riemannian four-manifold. There is an  $\epsilon_0$  so that for any proper subball  $U' \subset U$  the set of connections the restriction map*

$$M^{\epsilon_0}(U) \rightarrow M(U')$$

*is a compact map.*

Proof: Take a sequence of gauge equivalence classes of connections  $[A_i]$ . We can choose the representatives  $A_i = \Gamma + a_i$  so that they are in Coulomb gauge and we have uniform  $L_1^2$ -estimates on  $a_i$ . Pass to a subsequence where  $a_i$  converge weakly in  $L_1^2$ . By Fubini's theorem for each  $i$   $a_i|_{\partial B_r}$  is in  $L_1^2$  except for  $r$  in a set of measure zero. Since the sequence is countable for all  $i$ ,  $a_i|_{\partial B_r}$  a.e in  $r$ . Now fix such an  $r$ . Since the sequence is uniformly bounded we can pass to a subsequence where  $a_i|_{\partial B_r}$  is uniformly bounded in  $L_1^2$ . Now we can pass to a subsequence where  $a_i|_{\partial B_r}$  converges strongly in  $L_{1/2}^2$ . In particular  $CS(A_i|_{\partial B_r})$  converges. We now longer have that  $a_i|_{B_r}$  is in Coulomb gauge since it doesn't satisfy the boundary condition. However by the Big Slices Lemma we can put these  $a_i$  in to Coulomb gauge by a sequence of  $L_{5/2}^2$ -gauge transformations giving new connection one forms  $\tilde{a}_i$ .



We can assume the gauge transformations converge strongly in  $L^2_2 \cap C^0$  in particular  $\tilde{a}_i|_{\partial B_r}$  still converges in  $L^2_{1/2}$  and  $\tilde{a}_i$  still converges weakly in  $L^2_1$  to  $\tilde{a}$ . Set  $A = \Gamma + \tilde{a}$ . Since  $A_i$  are ASD, the convergence of CS implies  $\int_{B_r} |F_{A_i}|^2 * 1$  converges to  $\int_{B_r} |F_A|^2 * 1$ . This implies strong convergence of  $F_{A_i}$  and hence by the properness of the curvature map the strong convergence to  $\tilde{a}_i$  as required.  $\square$

We can also deduce regularity of solution from a similar discussion.

## Theorem

*Let  $A$  be an  $L^2_1$ -connection. Suppose that  $F_A^+ = 0$ . Then there is a an  $L^2_2$ -gauge transformations  $g$  so that  $gA$  is  $C^\infty$ .*

## Theorem

Let  $G$  be a compact Lie group. Let  $P \rightarrow X$  be principle  $G$ -bundle Let  $A_i$  be a sequence of ASD-connections with

$$\int_X |F_{A_i}|^2 \leq M$$

then after passing to a subsequence there are

so that for any open subset  $U \Subset X \setminus \{x_1, \dots, x_n\}$   $g_i \cdot A_i$  converges in the  $C^\infty$ -topology to  $A|_{X \setminus \{x_1, \dots, x_n\}}$ .

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- a bundle  $P' \rightarrow X$ .
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 $g_i : P|_{X \setminus \{x_1, \dots, x_n\}} \rightarrow P'|_{X \setminus \{x_1, \dots, x_n\}}$ ,

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- a connection  $A$  in a bundle  $P' \rightarrow X$

so that for any open subset  $U \Subset X \setminus \{x_1, \dots, x_n\}$   $g_i \cdot A_i$  converges in the  $C^\infty$ -topology to  $A|_{X \setminus \{x_1, \dots, x_n\}}$ .

Proof. We can pass to a subsequence where the sequence of curvature densities  $|F_{A_i}|^2 * 1$  converges in the weak topology to a density  $\omega$ . If we cover  $X$  by balls of radius  $2^{-N}$  then for each  $N$  so that the balls of half the radius cover as well. Then for each  $n$  we can find  $i(n)$  so that Then at most  $2M/\epsilon_0$  of these ball have

$$\int_B |F_{A_i}|^2 \leq \epsilon_0/2$$

for all  $i \geq i(N)$ . We can pass to a further subsequence so that the centers of these bad balls converge to points  $x_1, \dots, x_n$ . We thus get an exhaustion of  $X \setminus \{x_1, \dots, x_n\}$  by a sequence of balls  $B_j$  where

$$\int_{B_j} |F_{A_i}|^2 \leq \epsilon_0/2.$$

Finally by a diagonalization argument and the compact inclusion lemma we can pass to a subsequence where there are sequence  $g_n$  of gauge transformations to that

## Removeable singularities.

### Theorem (Uhlenbeck)

*Let  $P \rightarrow X$  be a principle bundle. Suppose that  $A$  is finite energy ASD connection on  $X \setminus \{x_1, \dots, x_n\}$ . Then there is a bundle  $P' \rightarrow X$  and a connection  $A'$  in  $P$  so that  $A|_{X \setminus \{x_1, \dots, x_n\}}$  and  $A'|_{X \setminus \{x_1, \dots, x_n\}}$  are gauge equivalent.*



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### Definition

A sequence of connection  $A_i$  in  $P \rightarrow X$  converges **upto bubbling** to  $A \in P' \rightarrow X$  if there is a collection of points  $\{x_1, \dots, x_n\}$  and a sequence of gauge transformations  $g_i$  so that for every compact subset  $K \Subset X \setminus \{x_1, \dots, x_n\}$  so that  $A_i|_K \mapsto A|_K$  in the  $C^\infty$ -topology.

Together with the compactness theorem this leads to the **Uhlenbeck compactification**.

### Definition

The Uhlenbeck compactification  $UM_k(X)$  of  $M_k(X)$  is the closure of  $M_k(X) \subset M_k(X) \cup X \times M_{k-1}(X) \cup \text{Sym}_2 X \times M_{k-1}(X) \cup \dots \cup \text{Sym}_k(X) \times M_0(X)$  with respect to the upto bubbling topology.

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### Remark

*The subset of  $M_k(X) \cup X \times M_{k-1}(X) \cup \text{Sym}_2 X \times M_{k-2}(X) \cup \dots \cup \text{Sym}_k(X) \times M_0(X)$  which actually appears in the closure is important to understand. Taubes' gluing work showed how to understand this subset in many situations.*

We will want to deal with orbifolds, orbifold bundles and orbifold connections. We'll only deal with  $\mathbb{Z}_2$ -orbifolds.  $X$  a Hausdorff space together with the following data. A collection of triples  $(\Gamma_i, V_i, \phi_i)$   $\Gamma_i$  is either  $\mathbb{Z}_2$  or the trivial group.  $V_i \subset \mathbb{R}^n$  is a  $\Gamma_i$  invariant open subset and  $\phi : V_i/\Gamma_i \rightarrow X$  is a homeomorphism onto an open subset  $U_i$ . The open subsets  $U_i$  are an open cover. Further for each inclusion  $U_i \hookrightarrow U_j$  there is an injective immersion  $\phi_{ij} : V_i \hookrightarrow V_j$  which intertwines the  $\Gamma_i$  and  $\Gamma_j$  actions. (Note that there is a unique injective homomorphism  $\Gamma_i \hookrightarrow \Gamma_j$  and we'd have to be a little more careful here if the  $\Gamma_i$  had non-trivial automorphisms.)

First for each  $(\Gamma_i, V_i, \phi_i)$  we have that  $\Gamma_i$  acts on  $V_i \times G$  commuting with the  $G$  action. A orbifold principal bundle is the data of transition functions

$$\tau_{ij} : V_i \times G \rightarrow V_j \times G$$

which again intertwines the  $\Gamma_i$  and  $\Gamma_j$  actions.

We will consider only very special 4-dimensional  $\mathbb{Z}_2$ -orbifolds. We allow only one local model where the  $\mathbb{Z}_2$  action is

$$(x_1, x_2, x_3, x_4) \mapsto (x_1, x_2, -x_3, -x_4).$$

Let  $F$  be the fixed point  $(x_1, x_2, 0, 0)$ . In this case the quotient  $\mathbb{R}^4/\mathbb{Z}_2$  has a natural smooth structure so we can think of  $X$  as smooth four manifold with a distinguished two dimensional submanifold  $S$ . We'll denote the orbifold by  $(X, S)$  though  $S$  might not globally be the fixed point set.  $S$  may or may not be orientable.

Suppose  $P$  is an orbifold  $SU_2$ -bundle and  $A$  is an orbifold connection in  $P$ . We get an bundle  $SU_2$  on  $X \setminus S$  with connectio.  $P$  may not extend over  $S$ . Suppose further that the  $\mathbb{Z}_2$  action on  $P$  is non trivial. Then over  $F$  we get map  $F$  to  $\text{Aut}(SU_2) = SO_3$ . Any element of order 2 in  $SO_3$  conjugate to an inner automorphism by the element

$$I = \begin{bmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{bmatrix}.$$

These inner automorphism form an  $\mathbb{RP}^2 \subset SO_3$ .

We can trivialize  $P$  so that the automorphism is locally constant Consider an invariant connection:

$$\Gamma + \begin{bmatrix} it_1 & z_1 \\ -\bar{z}_1 & -it_1 \end{bmatrix} dx^1 + \dots + \begin{bmatrix} it_4 & z_4 \\ -\bar{z}_4 & -it_4 \end{bmatrix} dx^4.$$

where  $t_1, t_2, z_3, z_4$  are even under the involution while  $t_3, t_4, z_1, z_2$  are odd. In particular along  $F$  the connection restricts to

$$\begin{bmatrix} it_1 & 0 \\ 0 & -it_1 \end{bmatrix} dx^1 + \begin{bmatrix} it_2 & 0 \\ 0 & -it_2 \end{bmatrix} dx^2$$

Examples. 1. Let  $\mathbb{Z}_2$  act on  $S^4$  by compactifying the standard action above. The fixed set is  $S^2$ . The action lifts to an action on the total space  $P_{\pm} = S^7 \rightarrow S^4$  and leaves the standard instanton (anti-instanton) invariant.  $X = S^4/\mathbb{Z}_2$  is homeomorphic to  $S^4$ . The moduli space of  $\mathbb{Z}_2$  invariant instantons is

$$\mathbb{H}^3 = SO_{3,1}/SO_3.$$

The action of such a connection

$$\kappa = \frac{1}{8\pi^2} \int_X \text{tr}(F_A \wedge F_A) = \pm 1/2$$

So for we have

$$d = 3, \kappa = 1/2, S \cdot S = 0, \chi(S) = 2, \chi(X) + \sigma(X) = 2.$$



2. Recall that  $\mathbb{C}\mathbb{P}^2 / \text{conj} = S^4$  (a folk observation made precise independently by Arnold, Kuiper, and Massey). Consider the map  $f : \mathbb{C}^3 \rightarrow \text{Sym}_o(\mathbb{R}^3)$ .

$$f : \mathbb{C}^3 \rightarrow \mathbf{z} \in S^5 \rightarrow \Re(\mathbf{z}\bar{\mathbf{z}}^t - \frac{1}{3}|\mathbf{z}|^2).$$

This is the real part of the moment map for  $SU_3$  acting on  $\mathbb{C}^3$ . The image moment map consists of traceless hermitian matrices with exact 2 eigenvalues. If  $\mathbf{z} \in S^5$  then  $\mathbf{z}$  is an eigenvalue  $2/3$  while if  $\mathbf{w}$  is orthogonal to  $\mathbf{z}$  the matrix acts by  $-1/3$ . Thus the image of the moment map is exactly  $\mathbb{C}\mathbb{P}^2$ . Each of these authors show that image of the real part is a copy of  $S^4$ . Call this orbifold  $(S^4, \mathbb{R}\mathbb{P}^2_-)$ .

The Fubini-Study metric gives and ASD connection in  $\text{End}_0(T^*\mathbb{CP}^2)$ .

$$\rho_1(T^*\mathbb{CP}^2) = c_1^2 - 4c_2 = (9 - 12)H^2 = -3H^2$$

so  $k = 3/4$ . The dimension of this moduli space is

$$6 - 3(1 + 1) = 0$$

and eventually one can show this connection is unique upto gauge. The involution on  $\mathbb{CP}^2$  is an isometry so lifts uniquely to  $T^*\mathbb{CP}^2$ . Pushing this down to the

Note that

$$d = 0, \kappa = 3/8, S \cdot S = -2, \chi(S) = 1, \chi(X) + \sigma(X) = 2.$$

3. If we take the same example but reverse orientation we have  $\overline{\mathbb{C}\mathbb{P}^2}/\mathbb{Z}_2 = S^4$ . On  $\overline{\mathbb{C}\mathbb{P}^2}$  there is a (unique upto gauge) reducible ASD connection in  $\mathbb{R} \oplus E$  where  $E$  is tautological bundle thought of on  $\overline{\mathbb{C}\mathbb{P}^2}$ . Now  $p_1(\mathbb{R} \oplus E) = -c_2(\mathbb{C} \oplus E \oplus E^{-1}) = -1$  so  $k = 1/4$ . Thus this moduli has formal dimension

$$8(1/4) - 3(1) = -1.$$

The  $-1$  is accounted for by the stabilizer of the connection. There are two  $\mathbb{Z}_2$  actions on  $\mathbb{R} \oplus E$  actions covering conjugation on  $\overline{\mathbb{C}\mathbb{P}^2}$ , these are  $\pm$  conjugation on  $E$  and  $-1$  on  $\mathbb{R}$ . These however are gauge equivalent by multiplication by  $i$  in  $E$ . The resulting connection on  $S^4$  has

$$d = 0, \kappa = 1/8, S \cdot S = 2, \chi(S) = 1, \chi(X) + \sigma(X) = 2.$$

4. Next  $S^2 \times S^2$  with involution which is conjugation on each factor. The quotient is again  $S^4$ , the involution acts by  $-1$  on  $H^2(S^2 \times S^2)$ ,  $S = S^1 \times S^1$ . Let  $L$  be line bundle which is the tensor product of the tautological bundle on the first factor and its dual on the second  $\pi_1^*(\mathcal{O}(-1)) \times \pi_2^*(\mathcal{O}(1))$ . With the product metric and both factors of the same area the curvature of  $L$  is ASD. The bundle  $\mathbb{R} \oplus L$  then has a unique reducible ASD connection.  $p_1(\mathbb{R} \oplus L) = c_1^2(L) = -2$  so the dimension of the moduli space is

$$4 - 6 = -2$$

Since the connection is reducible the deformation complex splits into the trivial part which has cohomology

$$H^0 = H^2 = \mathbb{R}, \text{ and } H^1 = 0$$

account for all of the index. The  $L$  part of the complex has trivial cohomology. The involution acts by  $-1$  on  $H^0$  and  $1$  on  $H^2$ , thus corresponding orbifold connection on  $S^4/T^2$

Using excision one can see that the dimension of the moduli space is an affine function of  $\kappa$ ,  $\chi(X) + \sigma(X)$ ,  $S \cdot S$ , and  $\chi(S)$  whether or on  $S$  is orientable:

$$d = 8\kappa - \frac{3}{2}(\chi(X) + \sigma) + aS \cdot S + b\chi(F) + c$$

From the examples we see that  $a = 1/2$ ,  $b = 1$ ,  $c = 0$  so we have

### Theorem

*The formal dimension of the moduli space  $M_\kappa(X, F)$  is*

$$d = 8\kappa - \frac{3}{2}(\chi(X) + \sigma) + \frac{1}{2}S \cdot S + \chi(S).$$

Compactness theorem for orbifold connections. Since locally orbifold connections are just invariant connections there is nothing more to prove although when curvature concentrates along  $F$  there is only a drop  $\kappa$  drops in units of  $1/2$  rather than 1.

To describe the Uhlenbeck compactification on the pair  $(X, S)$  consider weighted collections of points

$$\text{Sym}_*(X, S) = \left\{ \sum_i n_i x_i \mid n_i \in \mathbb{Z} \text{ if } x_i \in X \setminus S \text{ and } n_i \in \frac{1}{2}\mathbb{Z} \text{ if } x_i \in S \right\}$$

Then for  $\kappa \in \frac{1}{2}\mathbb{Z}$  set

$$\text{Sym}_\kappa(X, S) = \left\{ \sum_i n_i x_i \in \text{Sym}_*(X, S) \mid \sum_i n_i = \kappa \right\}.$$

Thus

$$\text{Sym}_{\frac{1}{2}}(X, S) = S$$

$$\text{Sym}_1(X, S) = X \cup_F \text{Sym}_2(F)$$

$$\text{Sym}_{\frac{3}{2}}(X, S) = (X \times F) \cup_{F \times F} \text{Sym}_3(F)$$

$$\text{Sym}_2(X, S) = (\text{Sym}_2(X) \cup \text{Sym}_2(F) \times X \cup \text{Sym}_4(F)) / \sim$$

etc.

Thus we have a compactification which now looks like

### Definition

The Uhlenbeck compactification  $UM_{\kappa}(X, S)$  of  $M_{\kappa}(X, S)$  is the closure of

$$M_{\kappa}(X, S) \subset \bigcup_{\kappa' \in \frac{1}{2}\mathbb{Z}} \text{Sym}_{\kappa'}(X, S) \times M_{\kappa - \kappa'}(X, S).$$

with respect to the upto bubbling topology.



## Floer's Instanton Homology

Let  $Y$  be a 3-manifold  $P = Y \times SU_2$ .  $\Gamma$  the trivial connection and  $B = \Gamma + b$ .

$$CS(B) = \int_Y \text{tr}(b \wedge db + \frac{1}{3}b \wedge [b \wedge b])$$

Floer's great insight was that using the analytic tools developed by Uhlenbeck, Taubes and Donaldson one could carry over the Morse complex construction to CS viewed as a function on  $\mathcal{B} = \mathcal{A}/\mathcal{G}$ .

To compute the directional derivative  $\frac{d}{ds} CS(B + sc)|_{s=0}$  introduce the one parameter family of  $4d$ -connections on  $[0, 1] \times Y$  with  $t$  being the coordinate in  $[0, 1]$

$$A_s = B + stc$$

Then

$$F_{A_s} = F_B + d_B^{(4)}(stc) + \frac{s^2 t^2}{2} [c \wedge c].$$

so that

$$\begin{aligned} \frac{d}{ds} F_{A_s}|_{s=0} &= dt \wedge d_B c \\ \frac{d}{ds} CS(B + sc)|_{s=0} &= \frac{d}{ds} \int_{[0,1] \times Y} \text{tr}(F_{A_s} \wedge F_{A_s})|_{s=0} \\ &= 2 \int_0^1 \int_Y dt \wedge \text{tr}(c \wedge F_B) \\ &= 2 \int_Y \text{tr}(c \wedge F_B) \end{aligned}$$

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Thus critical points of CS are flat connections  $F = 0$

The gauge group  $\mathcal{G} = \{g : Y \rightarrow SU_2\}$  acts on  $\Omega^1(Y) \otimes su_2$  by

$$g \cdot b = gbg^{-1} + gdg^{-1}.$$

Under this action we have

$$CS(g \cdot B) = CS(B) + \deg(g) \quad \text{and} \quad F_{g \cdot B} = gF_Bg^{-1}$$

Thus the set of critical points of  $CS$  is preserved by the  $\mathcal{G}$  action and

$$\{a | F_B = 0\} = \text{Hom}(\pi_1(Y), SU_2) / \text{conj} = R(Y)$$

where the identification is by the holonomy representation.

## Examples:

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- (Fintushuel-Stern)  $S$  a Seifert fibered homology sphere with **four** exceptional fibers. The representation space is a union of  $S^2$ 's

Given a riemannian metric on  $Y$  we can formally compute the gradient of  $CS$  and we have

$$2 \int_Y \text{tr}(F_B \wedge c) = 2 \int_Y \text{tr}(*F_B \wedge *c) = -2 \langle *F_B, c \rangle.$$

The downward gradient flow for  $\mathcal{L} = -\frac{1}{2}CS$  is

$$\frac{B}{dt} = -\nabla_B \mathcal{L} = \frac{1}{2} \nabla_B CS = - * F_B.$$

A **major miracle** occurs here. If we view a solution  $B(t)$  as a connection  $A$  on  $\mathbb{R} \times Y$  the connection satisfies the Anti-Self-Dual Yang-Mills equation

$$F_A = - * _4 F_A.$$

As a 4d-connection

$$F_A = dt \wedge \frac{dB}{dt} + F_B$$

and

$$*F_A = dt \wedge * _3 F_B + * _3 \frac{dB}{dt}$$

This later equations is invariant not just under  $\mathcal{G}_3 = \{g : Y \rightarrow SU_2\}$  but under  $\mathcal{G}_4 = \{h : \mathbb{R} \times Y \rightarrow SU_2\}$ .

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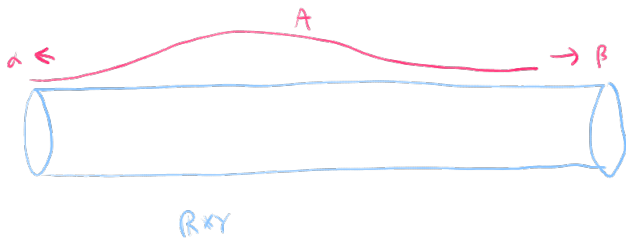


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- $\mathcal{B}$  is singular.
- $CS$  is not single valued (cf Novikov homology).
- Compactification of moduli spaces (including gluing).
- Construct perturbations to achieve Morse-Smale Condition that do not destroy compactness properties.

Given a pair of critical points  $\alpha, \beta$  set

$$M(\alpha, \beta) = \left\{ A \mid - \int_Z \text{tr}(F_A \wedge *F_A) < \infty, F_A = - * F_A, \right. \\ \left. \lim_{t \rightarrow -\infty} [A|_{t \times Y}] = \alpha, \lim_{t \rightarrow \infty} [A|_{t \times Y}] = \beta \right\} / \mathcal{G}_4.$$



Note that  $A$  is ASD if and only if

$$-\int_Z \text{tr}(F_A \wedge *F_A) = -CS(\alpha) + CS(\beta)$$

where the difference is computed with respect the path  $B(t) = A|_{t \times Y}$ .

Fix  $Q \rightarrow Y$  an  $SO_3$  or  $SU_2$  bundle. Given  $B_{\pm}$  flat connections **with trivial stabilizer** choose a path  $B(t)$  with  $t \in \mathbb{R}$  so that for  $t < -1$   $B(t) = B_-$  and  $t > 1$   $B(t) = B_+$ . Let  $A$  be the corresponding 4d-connection on and  $\mathbb{R} \times Q = P \rightarrow Z = \mathbb{R} \times Y$ . Then we make the following definitions.

$$\mathcal{A}(\alpha, \beta) = A + L_2^2(Z; T^*Z \otimes \text{ad}P)$$

$$\mathcal{G} = g \in 1 + L_3^2(Z, \text{End}(E)) | gg^* = 1.$$

Since  $gA$  is locally  $gag^{-1} + gag^{-1}$   $\mathcal{G}$  acts on  $\mathcal{A}(\alpha, \beta)$

$$\mathcal{B}(\alpha, \beta) = \mathcal{A}(\alpha, \beta) / \mathcal{G}.$$

## Remark

*If stabilizer of end points is not  $Z(G)$  this is not a good definition.*

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*The definition depends on a homotopy class  $\gamma \in \pi_1(\mathcal{B}_Q; \alpha, \beta)$  so write  $\mathcal{A}_\gamma(\alpha, \beta)$  or  $\mathcal{B}_\gamma(\alpha, \beta)$  when necessary.*

For  $A \in \mathcal{A}(\alpha, \beta)$

$$F_A^+ \in L_1^2(Z, \Lambda^+(Z) \otimes \text{ad}P)$$

Set

$$M_\gamma(\alpha, \beta) = \{[A] \in \mathcal{B}(\alpha, \beta) \mid F_A^+ = 0\}$$

We'd like to see that  $M_\gamma(\alpha, \beta)$  has a Kuranishi local structure and that

$$M(\alpha, \beta) = \bigcup_{\gamma \in \pi_1(\mathcal{B}_Q, \alpha, \beta)} M_\gamma(\alpha, \beta)$$

Write  $A = B + b + cdt$  where  $B$  is a pull back connection.  
Then

$$F_A = F_B + dt \wedge \dot{b} + d_B b - dt \wedge d_B c + \frac{1}{2}[b \wedge b] + dt \wedge [c \wedge b]$$

$$0 = F_A^+ = dt \wedge (\dot{b} + *F_B + d_B c + *d_B b + *\frac{1}{2}[b \wedge b] + [c \wedge b]) \\ + *\dot{b} + F_B + *d_B c + d_B b + \frac{1}{2}[b \wedge b] + *[c \wedge b].$$

The slice condition

$$0 = -\mathbf{d}_B^*(b + cdt) = *_4 \mathbf{d}_B *_4 (b + cdt) \\ = *\mathbf{d}_B dt \wedge *b + *c \\ = *_4(-dt \wedge (d_B *b + dt \wedge *\dot{c} + d_B *c)) \\ = \dot{c} + d_B^* b + dt \wedge d_B^* c.$$



Linearized equations at  $B$ .

$$0 = \begin{bmatrix} \dot{b} \\ \dot{c} \end{bmatrix} + \begin{bmatrix} *d_B & d_B \\ d_B^* & 0 \end{bmatrix} \begin{bmatrix} b \\ c \end{bmatrix}$$

This of the form  $\frac{d}{dt} + D$  where  $D$  is a first order self-adjoint elliptic operator.

### Theorem

*If  $D$  is an invertible first order self-adjoint elliptic operator acting on a vector bundle  $F \rightarrow Y$  then*

$$\frac{d}{dt} + D : L_k^2(Z; F) \rightarrow L_{k-1}^2(Z, F)$$

*is invertible.*

Sketch of proof in the case  $k = 1$ . Note that

$$\begin{aligned} \left\| \left( \frac{d}{dt} + D \right) u \right\|_{L^2(Z)}^2 &= \int_Z \left\langle \left( \frac{d}{dt} + D \right) u, \left( \frac{d}{dt} + D \right) u \right\rangle * 1 \\ &= \int_Z \left| \frac{du}{dt} \right|^2 + |Du|^2 + 2 \left\langle \frac{du}{dt}, Du \right\rangle * 1 \\ &= \int_Z \left| \frac{du}{dt} \right|^2 + |Du|^2 + \frac{d}{dt} \langle u, Du \rangle * 1 \\ &= \int_Z \left| \frac{du}{dt} \right|^2 + |Du|^2 * 1 \\ &\geq C \|u\|_{L^2_1}^2. \end{aligned}$$

Thus  $\frac{d}{dt} + D : L^2_k(Z; F) \rightarrow L^2_{k-1}(Z, F)$  is injective with closed range. Exercise: Show the range is dense. Generalize this to the case of general  $k$ .  $\square$ .

This can be used to prove in general that provided the Extend Hessian of  $\mathcal{L}$

$$EH_B = \begin{bmatrix} *d_B & d_B \\ d_B^* & 0 \end{bmatrix}$$

is invertible for  $B = B_{\pm}$  the linearized ASD equations and gauge fixing are Fredholm so the package we use to analyze the moduli space on closed manifolds carries over to this case.  $EH_B$  plays the role of the Hessian in finite dimensional (and no group action) Morse theory.

What does invertibility mean? If

$$\begin{bmatrix} b \\ c \end{bmatrix}$$

is in the kernel of  $EH_B$  and  $B$  is flat the  $b$  is a harmonic representative of  $H^1(Y, \text{ad}_B)$  and  $c \in H^0(Y, \text{ad}_B)$ . In other words the  $B$  must be

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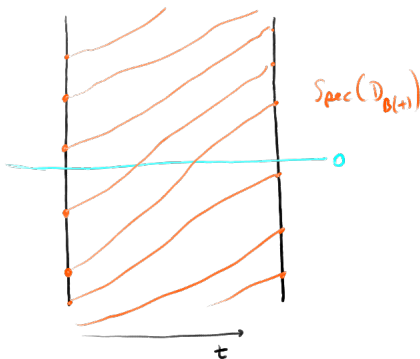
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- Irreducible
- Infinitesimally isolated (non-degenerate critical point.)

Given  $A \in \mathcal{A}_\gamma(\alpha, \beta)$  write  $A = B + cdt$  where  $B = B(t)$  is path in  $\mathcal{A}_Q$  and we can ask what is the index of

$$\frac{d}{dt} + D_B.$$

We now have a family of self-adjoint operators  $D_{B(t)}$ . Such a family has a **spectral flow**. Below we have a spectral



flow of  $+2$ .

Note the self-adjoint Fredholm operators  $\mathcal{SF}$  has three components

$$\mathcal{SF} = \mathcal{SF}_- \cup \mathcal{SF}_0 \cup \mathcal{SF}_+$$

where  $\mathcal{SF}_\pm$  are the essentially positive and negative operators.  $\mathcal{SF}_\pm$  are contractible while

$$\Omega\mathcal{SF}_0 \equiv \mathbb{Z} \times BO \text{ (or } \mathbb{Z} \times BU \text{ for complex ops).}$$

Theorem (Atiyah-Patodi-Singer..)

$$\text{Ind}\left(\frac{d}{dt} + D_t\right) = \text{sf}(D_t)$$

Idea of proof. Index is homotopy invariant so homotope the family so that  $D_t$  all have the same eigenvectors. Then consider

$$T_{\pm} = \frac{d}{dt} + \pm \tanh(t) : L^2_1(\mathbb{R}) \rightarrow L^2(\mathbb{R}).$$

The kernel of  $T_+$  is spanned by  $\operatorname{sech}(t)$  while for  $T_-$  the kernel is spanned by  $\cosh(t)$ . Since  $T_{\pm}^* = -T_{\mp}$  when

$$\operatorname{Ind}(T_{\pm}) = \pm 1 = \operatorname{sf}(\pm \tanh(t)).$$

□.



Note the spectral flow is the what remains of the Morse index.  $\mu(\alpha)$  should be the number of negative eigenvalue of the Hessian. Infinite in this case. However in the finite dimensional situation

$$\dim(M(x, y)) = \mu(y) - \mu(x) = \mathit{sf}(Hess_{\gamma(t)} f)$$

and we can make sense in this case of the difference

$$\mu(\beta) - \mu(\alpha) = \mathit{sf}(EH_{\gamma(t)})$$

If  $\gamma : S^1 \rightarrow \mathcal{B}_Q$  is a **closed** loop i.e. or  $\tilde{\gamma} : [0, 1] \rightarrow \mathcal{A}_Q$  and so that there is a gauge transformation with  $g\tilde{\gamma}(0) = \tilde{\gamma}(1)$  then  $sf$  may be non-zero! Indeed let  $A$  be the connection in  $S^1 \times Q$  with

$$(D_\gamma(t)) = \text{Ind}(D_A) = 8k - \frac{3}{2}(0 + 0) = 8k = 8 \deg(g).$$

N.B. if  $Q$  is an  $SO_3$  bundle  $k \in \frac{1}{2}\mathbb{Z}$ .